# Algorithms for Big Data (XIII) 

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## Review

We studied random walks on general graphs using spectral decomposition.
We introduced the notion of electrical networks.
We derived bounds on the cover time of random walks.

## Electrical Network

Now we formally justify the electrical network argument used last week.
For an edge with weight $\mathcal{w}_{e}$, we define its resistance $\mathrm{r}_{e}=\mathcal{w}_{e}^{-1}$.
For an edge $\{u, v\}$, we can assign numbers $\mathbf{i}(u, v)=-\mathbf{i}(v, u)$ as the current on the edge.
The collection of currents is required to satisfy Kirchhoff's law.
Ohm's law is used to define the potential drop between two ends of an edge.

## Matrix Form

It is instructive to express physical laws in the matrix form.
We use an ordered pair $(u, v)$ satisfying $u \leq v$ to represent an edge $\{u, v\} \in E$. The signed edge-vertex adjacency matrix $\mathrm{U} \in\{0,1,-1\}^{\mathrm{E} \times V}$ is defined as

$$
U((u, v), w)= \begin{cases}1 & \text { if } w=u \\ -1 & \text { if } w=v \\ 0 & \text { otherwise }\end{cases}
$$

Let $W \in \mathbb{R}^{\mathrm{E} \times \mathrm{E}}$ be $\operatorname{diag}\left(w\left(e_{1}\right), \ldots, w\left(e_{|\mathrm{E}|}\right)\right)$.

We use $\mathbf{i} \in \mathbb{R}^{\mathrm{E}}$ to denote the vector of currents, $\mathbf{v} \in \mathbb{R}^{V}$ to denote the vector of voltages.
It holds that

$$
\mathbf{i}=\mathrm{W} \cdot \mathrm{U} \cdot \mathbf{v}
$$

We use $\mathbf{i}_{\text {ext }}(u)$ to denote the amount of current entering $u$ externally.
Then $\mathbf{i}_{\text {ext }}(u)=\sum_{v \in N(u)} \mathbf{i}(u, v)$, and

$$
\mathbf{i}_{\mathrm{ext}}=\mathrm{U}^{\mathrm{T}} \mathbf{i}=\mathrm{U}^{\mathrm{T}} \cdot \mathrm{~W} \cdot \mathrm{U} \cdot \mathbf{v}
$$

If $\mathbf{i}_{\text {ext }}(u)=0$, we call it a internal node, otherwise, we call it a boundary node.

## Graph Laplacian

The matrix $\mathrm{L} \triangleq \mathrm{U}^{\mathrm{T}} \mathrm{WU}$ is again graph Laplacian.
Consider the spectral decomposition of L:

$$
\mathrm{L}=\sum_{i>1} \lambda_{i} \mathbf{v}_{i} \mathbf{v}_{i}^{\mathrm{T}}
$$

Using the decomposition, the equation becomes to

$$
\sum_{i \geq 1} a_{i} \mathbf{v}_{i}=\left(\sum_{i>1} \lambda_{i} \mathbf{v}_{i} \mathbf{v}_{i}^{\mathrm{T}}\right)\left(\sum_{i \geq 1} b_{i} \mathbf{v}_{i}\right)
$$

where $\mathbf{i}_{\text {ext }}=\sum_{i \geq 1} a_{i} \mathbf{v}_{i}$ and $\mathbf{v}=\sum_{i \geq 1} b_{i} \mathbf{v}_{i}$.

Therefore, we must have $a_{1}=0$, which means the current entering the network is equal to the current leaving the network!

Define the Moore-Penrose pseudo-inverse of L

$$
L^{+}=\sum_{i>1} \lambda_{i}^{-1} \mathbf{v}_{i} \mathbf{v}_{i}^{\mathrm{T}}
$$

Given $\mathbf{i}_{\text {ext }}$, we can compute $\mathbf{v}$ as long as we can compute $\mathrm{L}^{+}$.
We shift $\mathbf{v}$ so that

$$
\mathbf{v}=\mathrm{L}^{+} \mathbf{i}_{\mathrm{ext}} .
$$

## Effective Resistance

We are now able to formally define effective resistance.

$$
\mathrm{R}_{\mathrm{eff}}(u, v) \triangleq\left(\mathbf{e}_{u}-\mathbf{e}_{v}\right)^{\mathrm{T}} \mathrm{~L}^{+}\left(\mathbf{e}_{u}-\mathbf{e}_{v}\right)
$$

To see this, assuming one unit of current enters $u$ and leaves $v$ :

$$
\mathbf{v}=\mathrm{L}^{+}\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)
$$

On the other hand,

$$
\mathbf{v}(u)-\mathbf{v}(v)=\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)^{\mathrm{T}} \mathbf{v}=\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)^{\mathrm{T}} \mathrm{~L}^{+}\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right) .
$$

Note that $\mathrm{L}^{+}$is positive semi-definite, we can define

$$
L^{+/ 2}=\sum_{i>1} \lambda^{-1 / 2} \mathbf{v}_{i}
$$

Then we can write

$$
\mathbf{v}(\mathfrak{u})-\mathbf{v}(v)=\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)^{\mathrm{T}} \mathbf{v}=\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)^{\mathrm{T}} \mathrm{~L}^{+}\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)=\left\|\mathrm{L}^{+/ 2}\left(\mathbf{e}_{\mathfrak{u}}-\mathbf{e}_{v}\right)\right\|_{2}^{2}
$$

Examples: Series and Parallel graphs.

## Approximating Effective Resistance

Directly computing effective resistance requires to compute $L^{+}$, which is costly.
We can view $\mathrm{L}^{+/ 2} \mathbf{e}_{\mathrm{u}}$ and $\mathrm{L}^{+/ 2} \mathbf{e}_{v}$ as two vectors in $\mathbb{R}^{n}$ and approximate their distance using metric embedding technique.

Recall in Lecture 6, we learnt:

## Theorem

For any $0<\varepsilon<\frac{1}{2}$ and any positive integer $m$, consider a set of $m$ points $S \subseteq \mathbb{R}^{n}$. There exists an matrix $A \in \mathbb{R}^{k \times n}$ where $k=O\left(\varepsilon^{-2} \log m\right)$ satisfying

$$
\forall \mathbf{x}, \mathbf{y} \in S, \quad(1-\varepsilon)\|\mathbf{x}-\mathbf{y}\| \leq\|A \mathbf{x}-A \mathbf{y}\| \leq(1+\varepsilon)\|\mathbf{x}-\mathbf{y}\| .
$$

In our proof of JLT, each entry of the matrix $\mathcal{A}$ is from $\mathcal{N}(0,1 / k)$.
We only need to show how to compute $A L^{+/ 2}$ efficiently...
Let $\mathrm{L}^{\prime} \triangleq \mathrm{W}^{1 / 2} \mathrm{U}$, then $\left(\mathrm{L}^{\prime}\right)^{\mathrm{T}} \mathrm{L}^{\prime}=\mathrm{L}^{+}$.
Therefore $\left\|L^{\prime}\left(\mathbf{e}_{u}-\mathbf{e}_{v}\right)\right\|_{2}^{2}=R_{\text {eff }}(u, v)$.
We only need to solve d-linear equations in $L$ to obtain $A L^{\prime} L$.

